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Wave propagation in cross-linked random fiber networks

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We numerically investigate the propagation of small-amplitude elastic waves in random fiber networks. Our analysis reveals that the dynamic response of the system is not only controlled by its overall elasticity, but also by the local microstructure. In fact, we find that the longest fiber-segment plays a key role in dynamics when the network is excited with waves of short wavelength. In this case, the Bloch modes are highly non-affine as the longest segments oscillate close to their resonances. Based on this observation, we predict the low frequency dispersion curves of random fiber networks. © 2015 AIP Publishing LLC. [<http://dx.doi.org/10.1063/1.4936327>]

Disordered fibrous networks are ubiquitous and can be found in sound and thermal insulation systems,¹ paper and paper products,² filtration and liquid absorption systems,³ biological membranes and connective tissues,^{4–6} just to name a few. The structural behavior of these networks often plays a key role, motivating a deep interest in their mechanical response.⁷ It is now well recognized that the mechanics of fibrous networks depends not only on the elastic properties of the constituent fibers but also on the network connectivity.^{1,2} Furthermore, it has been shown that the deformation transitions from affine to non-affine^{8–10} as the network density or fiber cross-sectional area decreases.¹¹ However, despite the fact that fiber networks are often subjected to dynamic loadings, previous studies have primarily focused on the static behavior of random networks.^{12–16} Surprisingly, while the dynamic behavior of periodic beam networks has recently attracted considerable interest,^{17–20} the propagation of elastic waves in random fibrous networks has not yet been studied.

By contrast, the propagation of elastic waves in random spring networks^{23,24} and disordered packing of particles^{25–27} have been extensively studied. In particular, it has been shown that close to the jamming point (defined as the point where disordered matter transition from a flowing to a rigid state) such systems are anomalous in that they exhibit an excess of low frequency modes, and they do not appear to exhibit any ordinary continuum behavior.^{25–27} However, the response of random fiber networks, in which the bending deformation mode of fibers is taken into account (i.e., networks of beams), is significantly different, since such systems are rigid for average coordination numbers, \bar{z} , below the jamming point (i.e., for $\bar{z} < 4$).²⁶

In this letter, we investigate numerically the propagation of small-amplitude elastic waves in elastically stable random fibrous networks with $3 < \bar{z} < 4$, focusing on the first and second modes. We consider two-dimensional random networks composed of cylindrical fibers of length L_0 , and radius r made of an isotropic, linearly elastic material with Young's modulus E_f , shear modulus G_f , and density ρ_f . The networks are generated by depositing randomly N fibers of equal length L_0 in a periodic unit cell of dimensions $a \times a$. Cross-links (welded joints) are then introduced at all points where

the fibers intersect, resulting in a Poisson distribution for the lengths of fiber-segments.^{28,29}

It has been shown that the key parameters governing the static response of such fiber networks are the network density, $\rho = NL_0/a^2$, and the ratio between bending and axial stiffness of fibers, $l_b = \sqrt{(E_f I_f)/(E_f A_f)} = r/2$ (where I_f and A_f denote the moment of inertia and the cross-sectional area of the fibers, respectively).^{11,28} We choose the non-dimensional system parameters $\bar{\rho} = \rho L_0$, $\bar{r} = r/L_0 = 2l_b/L_0$, $\bar{\lambda} = \lambda/L_0$, and $\bar{\omega} = \omega \bar{r}/\omega_0$ to present the results, where λ and ω are the wavelength and frequency of the propagating elastic waves, respectively. Moreover, ω_0 is the first resonance frequency of a fiber-segment of length L_0 and radius r , assuming clamped boundary conditions at both ends (i.e., $\omega_0 = 22.37 \sqrt{(E_f r^2)/(4\rho_f L_0^4)}$ ³⁰). Finally, since we will find that the dynamic response of the networks is controlled also by the length of their longest fiber-segment, L_{max} , we introduce the dimensionless parameter $\bar{L}_{max} = L_{max}/L_0$.

To study the propagation of small-amplitude elastic waves in such networks, we perform frequency-domain wave propagation analysis within the finite element (FE) framework using the commercial package Abaqus/Standard. Models of the periodic unit cell are built using Euler-Bernoulli beam elements (Abaqus element type B23), and Bloch-type boundary conditions are applied to the edges of the unit cell^{31,32} (see Ref. 33 for details on the wave propagation analysis). We then calculate the dispersion relation $\omega = \omega(\mathbf{k})$ using the frequency-domain perturbation method.

We start by computing the density of vibrational states for a network with $\bar{\rho} = 25$ and $\bar{r} = 0.004$ (see Fig. 1(c)). The results reported in Fig. 1(a) are based on a total of 2500 \mathbf{k} -space points uniformly distributed across the entire first Brillouin zone and are normalized by the number of \mathbf{k} -space points. In stark contrast with the case of a network of springs, whose density of state plot is dominated by a peak at $\omega \sim 0$ (see Fig. S2 in Ref. 33), the response of the fiber network is characterized by several peaks at non-zero frequencies. The first peak occurs at $\bar{\omega} = 0.011$, which corresponds to the resonance frequency of the longest fiber-segment $\bar{L}_{max} = 0.59$.

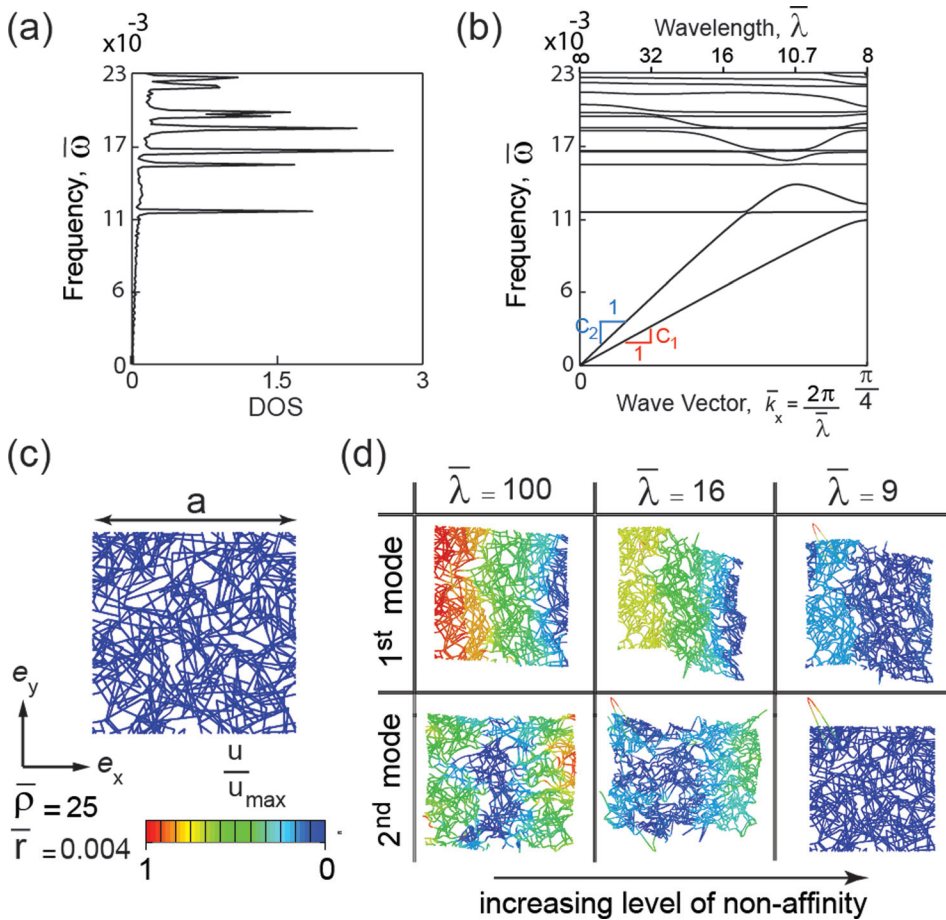


FIG. 1. (a) Density of vibrational states and (b) dispersion curves of a 2D cross-linked random network with $\bar{\rho} = 25$ and $\bar{r} = 0.004$ shown in (c). (d) The first (shear-like) and second (dilatation-like) Bloch modes for three different wavelengths of the propagating waves. Note that the hybridization of the folded shear-like mode with dilation-like mode observed at $\bar{\omega} = 15$ in (b) (an artifact of Bloch-type boundary conditions) has already been observed for surface acoustic waves.^{21,22}

To gain more insight into the dynamic behavior of the network, in Fig. 1(b), we plot its dispersion band diagram in an arbitrary direction (here chosen to be the horizontal direction \mathbf{e}_x). Similar to the case of an homogeneous linear elastic material,³⁴ two bands emanate from the origin, corresponding to shear and dilatational waves propagating with speeds c_1 and c_2 in the long wavelength limit. However, as the wavelength decreases, the bands eventually reach the resonance frequency of the longest fiber-segment and flatten.³⁵ It is important to note that this transition has a profound effect on the Bloch modes (see Fig. 1(d)). While in the long wavelength limit ($\bar{\lambda} = 100$ in Fig. 1(d)), both the first and second modes are characterized by low degree of non-affinity, for wavelength shorter than that of the transition ($\bar{\lambda} = 9$ in Fig. 1(d)), the vibrations are fully localized within the longest fiber-segment (see the red fiber with maximum displacement in the second mode), visually suggesting that the level of non-affinity increases as the wavelength of the propagating waves decreases.³⁶

To study the dynamic response of the fiber network in detail, first we investigate the response of the system in the long wavelength regime and determine the speeds c_1 and c_2 associated with the first (transverse) and second (longitudinal) modes as a function of the microstructural parameters of the network. Data for 115 random fiber networks with network densities $\bar{\rho}$ ranging from 12.5 to 25, and \bar{r} ranging from 0.00004 to 0.08 are presented in Fig. 2. To collapse the data onto a master curve, the same scaling parameter as in the static analysis of the network elasticity is used for the horizontal axis (i.e., $w = \log_{10}[\bar{\rho}^7(\bar{r}/2)^2]$ ³⁷), and the

wave speeds are normalized by the corresponding bulk properties of the fiber material (i.e., $(c_f)_1 = \sqrt{G_f/\rho_f}$ and $(c_f)_2 = \sqrt{G_f(4G_f - E_f)/[\rho_f(3G_f - E_f)]}$ ³⁴). As for the static case, the master plot has two well-defined regions. It has been shown that the transition between those two regions correlates with the transformation of the deformation field of the networks from non-affine to affine, as well as with the shift of energy storage from bending to axial.³⁷ For small values of w , the slope of the master curve is 1/2, indicating that the wave speeds of the networks are proportional to $\bar{\rho}^{3.5}\bar{r}$ in the long wavelength regime. Thus, the

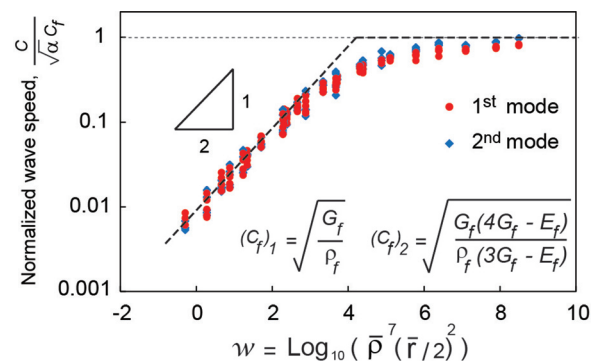


FIG. 2. Master plot providing the normalized wave speeds associated with the first and second modes as a function of system parameters. $\alpha = 0.38$ is a non-dimensional constant that has the same value as in the analysis of network elasticity.³⁷ $(c_f)_1$ and $(c_f)_2$ are the transverse and longitudinal wave speeds for the fibers material, respectively.

overall elastic modulus of networks, $E \sim c^2 \bar{\rho} r^2$,³⁸ is proportional to $\bar{\rho}^8 \bar{r}^4$, exactly as predicted by the static analysis. However, for $w > 6$, all data converge to a horizontal asymptote, so that the long wavelength speeds of the networks become independent of $\bar{\rho}$ and \bar{r} , and $E \sim \bar{\rho} \bar{r}^2$, again in full agreement with the static analysis.³⁷

Having determined the effect of the system parameters $\bar{\rho}$ and \bar{r} on the long wavelength response, we proceed to investigate the effect of the dimensionless wavelength, $\bar{\lambda}$, on the propagation of the elastic waves. To this end, we calculate the dispersion bands for the networks and also investigate the effect of $\bar{\lambda}$ on their Bloch modes by directly measuring the evolution of their degree of non-affinity. Note that the deviation from affine deformation, denoted by ΔA_j , for the j -th mode is defined as³⁹

$$\Delta A_j = \frac{1}{n U_0} \sum_{i=1}^n \|\mathbf{u}_j^i - \mathbf{u}_{aff,j}^i\|, \quad (1)$$

where n is the number of the nodes in a full sinusoidal period, U_0 is the amplitude of the propagating wave, \mathbf{u}_j^i is the displacement of the i -th node associated with the j -th mode, and $\mathbf{u}_{aff,j}^i$ the corresponding values when the deformation is fully affine. For a wave propagating along the \mathbf{e}_x direction, $\mathbf{u}_{aff,1}^i = U_0 \sin[2\pi(x_i - x_0)/\lambda] \mathbf{e}_y$, where x_i identifies the coordinate in \mathbf{e}_x direction of i -th node, \mathbf{e}_y is a unit vector

perpendicular to \mathbf{e}_x , and x_0 represents the phase of the mode.

In Figs. 3(a) and 3(b), we report the first dispersion band and the evolution of ΔA as a function of the normalized wavelength $\bar{\lambda}$ for two random networks characterized by the same values of $\bar{\rho}$ and \bar{r} ($\bar{\rho} = 25$ and $\bar{r} = 0.0004$), but different lengths of the longest fiber-segment ($\bar{L}_{max} = 0.55$ and 0.9).⁴⁰ Focusing on the dispersion bands (Fig. 3(a)), we see that the two networks are characterized by identical response in the long wavelength limit. We also find that in this region ΔA is not affected by $\bar{\lambda}$ (Fig. 3(b)). In fact, ΔA is the same as that calculated in the static case and is dictated by the overall elasticity of the system.^{13,28,41,42} By contrast, the wavelength $\bar{\lambda}^*$ and frequency $\bar{\omega}^*$, at which the band flattens is found to be significantly affected by \bar{L}_{max} . For the network with $\bar{L}_{max} = 0.55$ (red line), we find that $\bar{\lambda}^* = 22$ and $\bar{\omega}^* = 0.0012$, while if $\bar{L}_{max} = 0.9$ (black line), $\bar{\lambda}^* = 63$ and $\bar{\omega}^* = 0.0005$. We also note that for $\bar{\lambda} < \bar{\lambda}^*$, ΔA suddenly rises (see Fig. 3(b)), reflecting the emergence of fully localized modes (i.e., see the second mode for $\bar{\lambda} = 9$ in Fig. 1(d)).

Finally, the evolution of ΔA also reveals that the sharp transition at $\bar{\lambda}^*$ is preceded by another transition at $\bar{\lambda}'$, at which ΔA gradually increases (note that $\bar{\lambda}'$ is defined as the wavelength for which $\partial(\Delta A)/\partial \bar{\lambda} = 0.2\%$). Inspection of the modes reveals that this increase of ΔA is induced by the

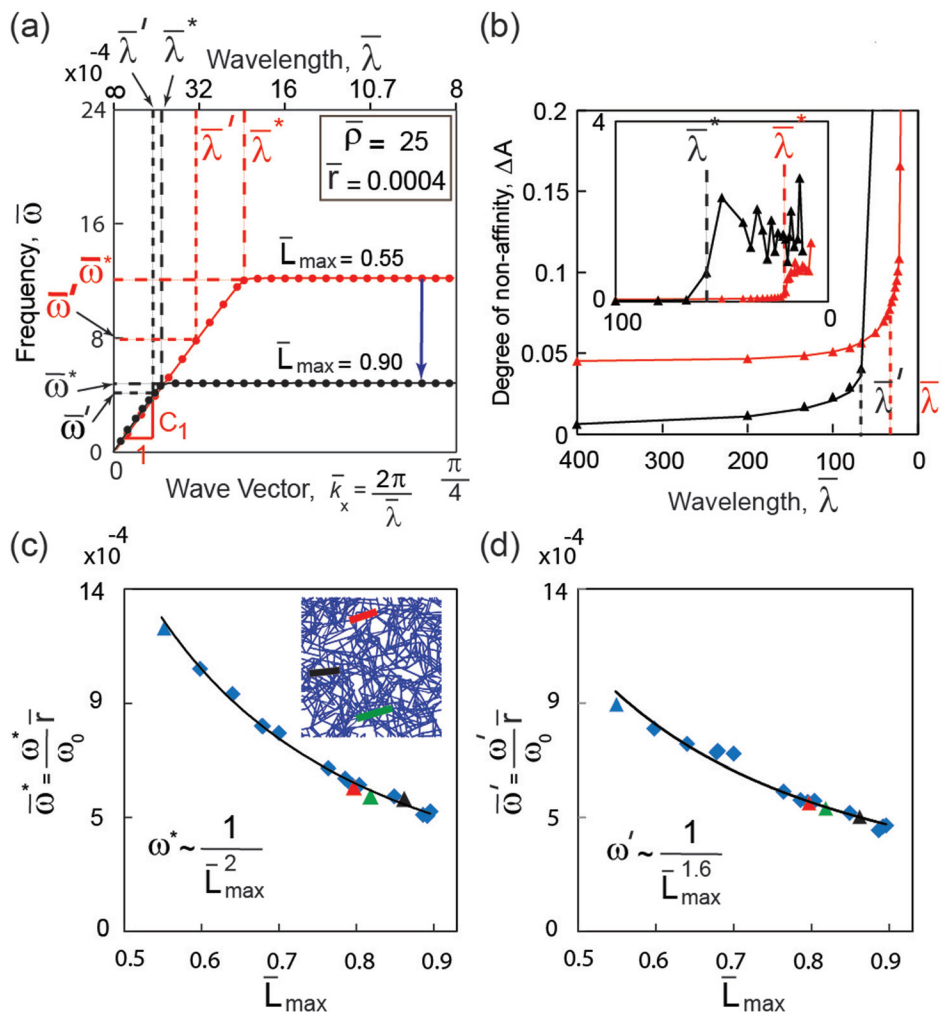


FIG. 3. (a) and (b) Effect of the longest fiber-segment, \bar{L}_{max} , on (a) dispersion curves, and (b) degree of non-affinity as a function of $\bar{\lambda}$ for networks with $\bar{\rho} = 25$ and $\bar{r} = 0.0004$. (c) and (d) Effect of \bar{L}_{max} on transition frequencies (c) $\bar{\omega}'$, and (d) $\bar{\omega}^*$, for different random networks (blue markers) with $\bar{\rho} = 25$, $\bar{r} = 0.0004$. The blue, red, green, and black triangular markers correspond to the exactly same network (shown as an inset in (c)), but with different \bar{L}_{max} attached to it (shown in the inset).

gradual localization of the deformation within the longest fiber-segment of the network (i.e., see the first and second modes for $\bar{\lambda} = 16$, and the first mode for $\bar{\lambda} = 9$ in Fig. 1(d)). Importantly, we find that in the vicinity of this transition, the dispersion curve starts deviating from linearity (see Fig. S3 in Ref. 33), so that $\bar{\lambda}'$ also identifies the threshold below which the fibrous network behaves as a dispersive wave medium. Interestingly, the results reported in Figs. 3(a) and 3(b) indicate that $\bar{\lambda}'$ also is significantly affected by L_{max} .

To fully understand the key role played by the longest fiber-segment in the dynamic response of random fiber networks, we calculate the characteristic frequencies, $\bar{\omega}' = \omega' \bar{r} / \omega_0$ and $\bar{\omega}^* = \omega^* \bar{r} / \omega_0$, for 16 different networks characterized by the same system parameters $\bar{\rho} = 25$ and $\bar{r} = 0.0004$, but different values of \bar{L}_{max} . The results reported in Fig. 3(c) show that $\bar{\omega}^*$ scales with $1/\bar{L}_{max}^2$. This scaling law derives directly from the expression for the resonance frequency of a fiber-segment of length L_{max} (i.e., $\omega^* = 22.37 \sqrt{(E_f I_f) / (\rho_f A_f L_{max}^4)}$), since this frequency corresponds to fully localization of vibrations within the longest fiber-segment. Moreover, we interpret the transition frequency $\bar{\omega}'$ as the onset of resonance of the longest fiber-segment, and expect that $\bar{\omega}' / \bar{\omega}^*$ is affected by the network environment surrounding the fiber of length L_{max} . In particular, we expect that $\bar{\omega}'$ is always smaller than $\bar{\omega}^*$, and $\bar{\omega}' / \bar{\omega}^* \rightarrow 1$ as \bar{L}_{max} increases (i.e., L_{max} becomes very large compared with the other segments of the network). Therefore, since $\bar{\omega}^*$ scales with $1/\bar{L}_{max}^2$, it is expected that $\bar{\omega}'$ scales with $1/\bar{L}_{max}^m$ in which $m < 2$. Fig. 3(d) shows $\bar{\omega}'$ is proportional to $1/\bar{L}_{max}^{1.6}$ for networks with $\bar{\rho} = 25$ and $\bar{r} = 0.0004$. The robustness of these dependencies is confirmed by the fact that when we attach a longer fiber of lengths $\bar{L}_{max} = 0.796$ (red triangular marker), $\bar{L}_{max} = 0.821$ (green triangular marker), or $\bar{L}_{max} = 0.861$ (black triangular marker) to the network with $\bar{L}_{max} = 0.55$ (blue triangular marker), $\bar{\omega}'$ and $\bar{\omega}^*$ are considerably altered, but are still proportional to $1/\bar{L}_{max}^{1.6}$ and $1/\bar{L}_{max}^2$, respectively.

So far, we focused on fiber networks with $\bar{\rho} = 25$ and $\bar{r} = 0.0004$, and showed that their long wavelength response is purely governed by these two parameters, while \bar{L}_{max} plays an important role as $\bar{\lambda}$ is reduced, and strongly affects the wavelength $\bar{\lambda}^*$ at which the vibrations localize within a single fiber. However, it is important to recognize that $\bar{\lambda}^*$ is not only controlled by \bar{L}_{max} , but also by $\bar{\rho}$ and \bar{r} . In fact, $\bar{\lambda}^*$ is determined by the intersection between the flat band at $\bar{\omega}^*$ (which is determined by \bar{r} and \bar{L}_{max}), and the band with wave speed c_1 in the long wavelength regime (which is controlled by $\bar{\rho}$ and \bar{r} , as predicted in Fig. 2).

To better understand these dependencies, in Fig. 4(a), we focus on networks with $\bar{r} = 0.0004$ and $0.55 < \bar{L}_{max} < 0.9$, and determine $\bar{\lambda}^*$ for $\bar{\rho} = 25, 20$, and 12.5 . Note that, for the sake of simplicity, $\bar{\lambda}^*$ is predicted without accounting for the deviation of the shear-dominated band from linearity occurring for $\bar{\lambda} < \bar{\lambda}'$, since this deviation does not alter the trends discussed here. All networks considered in Fig. 4(a) are characterized by the same range of possible values for $\bar{\omega}^*$ (grey shaded region), while c_1 increases as $\bar{\rho}$ increases (according to the master plot in Fig. 2). Consequently, the range of possible values for $\bar{\lambda}^*$ moves to longer wavelengths

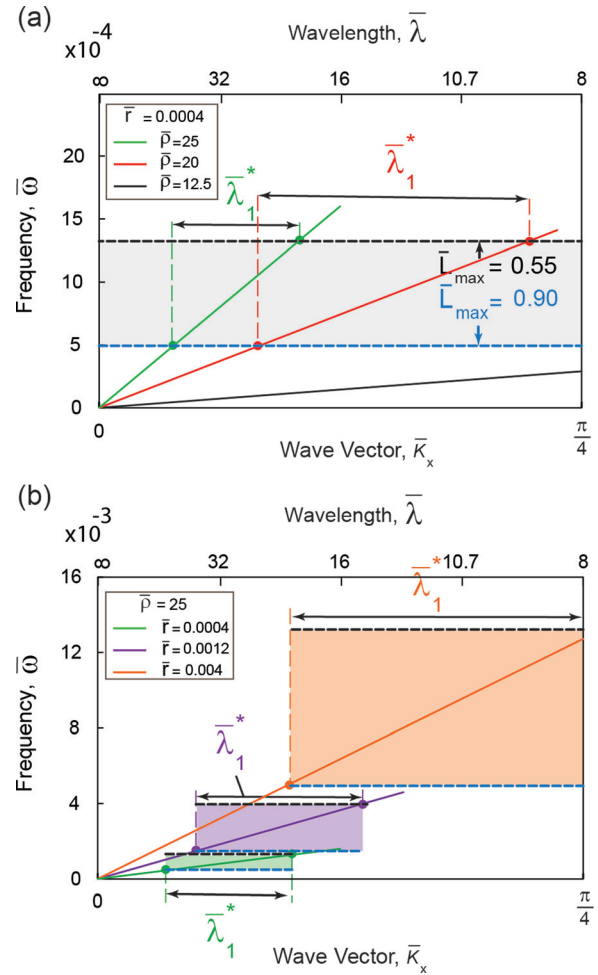


FIG. 4. Effects of (a) $\bar{\rho}$, and (b) \bar{r} on the wavelength $\bar{\lambda}^*$, at which the vibrations localize within a single fiber for the first mode. The shaded areas indicate the range of resonance frequencies for networks with $0.55 < \bar{L}_{max} < 0.9$.

as $\bar{\rho}$ increases. Note that such effect of $\bar{\rho}$ on $\bar{\lambda}^*$ is counter-intuitive. In fact, we would expect that for larger values of $\bar{\rho}$, waves of shorter wavelength are needed to trigger localized modes, since according to static behavior of random networks,¹¹ both overall stiffness and level of affinity of the system increase monotonically with $\bar{\rho}$.

On the other hand, if $\bar{\rho}$ remains constant and \bar{r} increases, the range of possible values for $\bar{\omega}^*$ increases quadratically, as illustrated in Fig. 4(b). Since c_1 increases at most linearly as \bar{r} increases (see master plot in Fig. 2), the slope of the shear band in the long wavelength regime increases slower than $\bar{\omega}^*$ limits. Hence, in networks with larger \bar{r} , highly localized modes are more likely triggered at shorter wavelengths, as illustrated in Fig. 4(b).

In summary, our analyses reveal that the dynamic response of the networks depends not only on their microstructural parameters, but also on the length of the longest fiber-segment and the wavelength of the propagating waves. Furthermore, we found that the propagation of the elastic waves through the system is characterized by three regimes: (i) for long wavelengths, the dynamic response of the system is fully controlled by the overall elasticity of the networks and is non-dispersive; (ii) for intermediate wavelengths, the networks behave as dispersive wave media and the non-affinity level of the Bloch modes increases

gradually; (iii) in the short wavelengths regime, the Bloch modes are highly non-affine with all deformation localized within the longest fiber-segment of the network. We also note that in this study, we focused on the dynamic response of 2D networks, while the study of more realistic 3D systems (whose behavior has been found to be qualitatively similar to that of 2D networks for the static case^{7,9}) will be the subject of future investigations. Our results not only shed light on the propagation of elastic waves in random fibrous networks, but also open avenues for the design of fiber networks for applications in which the structure is exposed to mechanical vibrations.

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³⁶Note that although in this study we investigate the dynamic response of 2D random fiber networks by focusing on a sufficiently large unit cell and applying Bloch-type boundary conditions, qualitatively the same behavior is found when we consider finite-size domains and calculate their steady-state dynamic linearized response to harmonic excitation (see Ref. 33 for details). In particular, also for finite-size domains, we find that localization happens when the frequency of the propagating waves is in the vicinity of the natural frequency of the longest fiber-segment, indicating that this is not an artifact introduced by Bloch-type boundary conditions (i.e., it is not caused by constructive scattering from neighboring unit cells).
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